

Further studies of isolated photon production with a jet in deep inelastic scattering at HERA

The ZEUS collaboration

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ABSTRACT: Isolated photons with high transverse energy have been studied in deep inelastic ep scattering with the ZEUS detector at HERA, using an integrated luminosity of 326 pb^{-1} in the range of exchanged-photon virtuality $10\text{--}350 \text{ GeV}^2$. Outgoing isolated photons with transverse energy $4 < E_T^\gamma < 15 \text{ GeV}$ and pseudorapidity $-0.7 < \eta^\gamma < 0.9$ were measured with accompanying jets having transverse energy and pseudorapidity $2.5 < E_T^{\text{jet}} < 35 \text{ GeV}$ and $-1.5 < \eta^{\text{jet}} < 1.8$, respectively. Differential cross sections are presented for the following variables: the fraction of the incoming photon energy and momentum that is transferred to the outgoing photon and the leading jet; the fraction of the incoming proton energy transferred to the photon and leading jet; the differences in azimuthal angle and pseudorapidity between the outgoing photon and the leading jet and between the outgoing photon and the scattered electron. Comparisons are made with theoretical predictions: a leading-logarithm Monte Carlo simulation, a next-to-leading-order QCD prediction, and a prediction using the k_T -factorisation approach.

KEYWORDS: Lepton-Nucleon Scattering (experiments), Photon production, QCD

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1 Introduction

The isolated high-energy photons that are emitted in high-energy collisions involving hadrons are predominantly unaffected by parton hadronisation. Their production probes the underlying partonic process and can provide information on the structure of the proton. Processes of this type have been studied in a number of fixed-target and hadron-collider experiments [1–10]. The production of isolated photons in photoproduction, where the incoming photon is quasi-real, was previously studied at HERA by the ZEUS and H1 collaborations [11–15]. Deep inelastic neutral current (NC) ep scattering (DIS), in which the exchanged photon has virtuality $Q^2 > 1 \text{ GeV}^2$, has also been measured in a variety of Q^2 ranges [16–18]. The analysis presented here extends an earlier ZEUS measurement of isolated photons and jets in DIS [19].

Figure 1 shows leading-order diagrams for high-energy photon production in DIS. Such “prompt” photons are emitted either by the incoming or outgoing quark or by the incoming or outgoing lepton. In the first case, the photons are classified as “QQ” photons, and the hadronic process has two hard scales: the virtuality Q^2 of the incident exchanged photon and the square of the transverse momentum of the prompt photon. In the second case, the photons are denoted as “LL” and are emitted from the incoming or outgoing lepton. The present analysis requires the observation of a scattered electron, a high-energy outgoing

photon and a hadronic jet. Processes in which the final state consists solely of a hard outgoing electron and a hard outgoing photon are thereby excluded. By requiring the outgoing photon to be isolated, a further class of processes in which the photon is produced within a jet is suppressed.

In the previous ZEUS publication on this topic [19], kinematic distributions of the outgoing photon and the jet were studied. Using the same data set, the analysis is extended here by measuring variables that involve two of the outgoing photon, the jet and the scattered electron. Results from a leading-logarithm parton-shower Monte Carlo [20] are compared to the measurements. Comparison is also made with two theoretical models: one at next-to-leading order (NLO) in QCD [21, 22], and one based on a k_T -factorisation approach [23].

2 Experimental set-up

The data sample used for the measurement corresponds to an integrated luminosity of $326 \pm 6 \text{ pb}^{-1}$ and was taken with the ZEUS detector in the years 2004–2007. During this period, HERA ran with an electron/positron beam energy of 27.5 GeV and a proton beam energy of 920 GeV; $138 \pm 2 \text{ pb}^{-1}$ of e^+p data and $188 \pm 3 \text{ pb}^{-1}$ of e^-p data¹ were used in the present analysis.

A detailed description of the ZEUS detector can be found elsewhere [24]. Charged particles were recorded in the central tracking detector (CTD) [25–27] and a silicon microvertex detector [28] which operated in a magnetic field of 1.43 T provided by a thin superconducting solenoid. The high-resolution uranium-scintillator calorimeter (CAL) [29–32] consisted of three parts: the forward (FCAL), the barrel (BCAL) and the rear (RCAL) calorimeters. The BCAL covered the pseudorapidity range -0.74 to 1.01 as seen from the nominal interaction point.² The FCAL and RCAL extended the range to -3.5 to 4.0 . The smallest subdivision of the CAL is called a cell. The barrel electromagnetic calorimeter (BEMC) cells had a pointing geometry aimed at the nominal interaction point, with a cross section approximately $5 \times 20 \text{ cm}^2$, with the finer granularity in the Z -direction. This fine granularity allows the use of shower-shape distributions to distinguish isolated photons from the products of neutral meson decays such as $\pi^0 \rightarrow \gamma\gamma$.

The luminosity was measured using the Bethe-Heitler reaction $ep \rightarrow e\gamma p$ by a luminosity detector which consisted of two independent systems: a lead-scintillator calorimeter [33–35] and a magnetic spectrometer [36].

3 Event selection and reconstruction

The ZEUS experiment operated a three-level trigger system [24, 37, 38]. At the first level, events were selected if they had an energy deposit in the CAL consistent with an isolated

¹Hereafter, “electron” refers to both electrons and positrons unless otherwise stated.

²The ZEUS coordinate system is a right-handed Cartesian system, with the Z axis pointing in the nominal proton beam direction, referred to as the “forward direction”, and the X axis pointing towards the centre of HERA. The coordinate origin is at the centre of the central tracking detector. The pseudorapidity is defined as $\eta = -\ln(\tan \frac{\theta}{2})$, where the polar angle, θ , is measured with respect to the Z axis. The azimuthal angle, ϕ , is measured with respect to the X axis.

electron. At the second level, a requirement on the energy and longitudinal momentum of the event was used to select NC DIS events. At the third level, the full event was reconstructed and tighter requirements for a DIS electron were made. Offline selections, similar to those of the earlier ZEUS analysis [19], were then applied.

Outgoing electrons were selected with polar angle $\theta_e > 140^\circ$ in order to provide a good measurement in the RCAL, kinematically separated from the selected outgoing photons. Their impact point (X, Y) on the surface of the RCAL was required to lie outside a rectangular region ± 14.8 cm in X and $[-14.6, +12.5]$ cm in Y , to give a well understood acceptance. The outgoing electrons were identified using a neural network [39], and the energy of the outgoing electron, E'_e , corrected for apparatus effects, was required to be larger than 10 GeV. The kinematic variable Q^2 was reconstructed as $Q^2 = -(k - k')^2$, where k (k') is the four-momentum of the incoming (outgoing) electron. The kinematic region $10 < Q^2 < 350$ GeV² was selected.

A requirement that the event vertex position, Z_{vtx} , should be within the range $|Z_{\text{vtx}}| < 40$ cm reduces the background from non- ep collisions. A further requirement for a well-contained DIS event, $35 < E - p_Z < 65$ GeV, was imposed where $E - p_Z = \sum_i E_i(1 - \cos \theta_i)$; E_i is the energy of the i -th CAL cell, θ_i is its polar angle and the sum runs over all cells [40].

Photon candidates were identified as energy-flow objects (EFOs)³ without an associated track, for which at least 90% of the reconstructed energy was deposited in the BEMC. The calibration of the energies of the photon and scattered electron was taken from an earlier ZEUS analysis and used deeply virtual Compton scattering events [44]. The reconstructed transverse energy of the photon candidate, E_T^γ , was required to lie within the range⁴ $4 < E_T^\gamma < 15$ GeV and the pseudorapidity, η^γ , had to satisfy $-0.7 < \eta^\gamma < 0.9$.

Jets were reconstructed with the k_T clustering algorithm [45] in the E scheme in the longitudinally invariant inclusive mode [46] with the R parameter set to 1.0. Since all EFOs of the event were used except for the electron signal, one of the jets found by this procedure corresponds to or includes the photon candidate. At least one accompanying jet was required with transverse energy $E_T^{\text{jet}} > 2.5$ GeV and pseudorapidity, η^{jet} , in the range $-1.5 < \eta^{\text{jet}} < 1.8$; if more than one jet was found, that with the highest E_T^{jet} was used.

Photons radiated from final-state electrons were suppressed by requiring that $\Delta R > 0.2$, where $\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2}$ is the distance to the nearest reconstructed track with momentum greater than 250 MeV in the $\eta - \phi$ plane. Isolation from hadronic activity was imposed by requiring that the photon candidate possessed at least 90% of the total energy of the jet-like object of which it formed a part. This also reduced the background of photon candidates arising from neutral meson decay.

Approximately 6000 events were selected at this stage; this sample was dominated by background events in which one or more neutral mesons such as π^0 and η , decaying to photons, produced a photon candidate in the BEMC.

³Energy-flow objects [41–43] were constructed from calorimeter-cell clusters and tracks, associated when possible.

⁴The upper limit was selected to retain distinguishable shower shapes between the hadronic background and the photon signal.

4 Variables studied

In the previous ZEUS publication [19], distributions of photon and jet variables were studied. In the present analysis, variables that depend on two of the three measured outgoing physical objects were studied, namely the high- p_T photon, the leading jet and the scattered electron. They were defined as follows:

- x_γ^{meas} is a measure of the fraction of the exchanged-photon energy and longitudinal momentum that is given to the outgoing photon and the jet:

$$x_\gamma^{\text{meas}} = \frac{E^\gamma - p_Z^\gamma + E^{\text{jet}} - p_Z^{\text{jet}}}{2E_e y_{\text{JB}}},$$

where E^γ and E^{jet} denote the energies of the outgoing photon and the jet, respectively, p_Z^γ and p_Z^{jet} denote the corresponding longitudinal momenta, $E_e = 27.5$ GeV, and the Jacquet-Blondel variable y_{JB} is given by $\sum_{\text{EFO}}(E^{\text{EFO}} - p_Z^{\text{EFO}})/2E_e$, summing over all energy-flow objects in the event except the scattered electron, each object being treated as equivalent to a massless particle. This variable is sensitive to higher-order processes that generate additional particles in the event;

- x_p^{obs} estimates the fraction of the proton energy transferred to the outgoing photon and jet:

$$x_p^{\text{obs}} = \frac{E^\gamma + p_Z^\gamma + E^{\text{jet}} + p_Z^{\text{jet}}}{2E_p},$$

where $E_p = 920$ GeV. This variable is sensitive to the partonic structure of the proton;

- $\Delta\phi$ is the azimuthal angle between the jet and the outgoing photon: $\Delta\phi = |\phi^{\text{jet}} - \phi^\gamma|$, where ϕ^{jet} and ϕ^γ denote the azimuthal angles of the jet and photon, respectively. This variable is sensitive to the presence of higher-order gluon radiation from the outgoing quark, which generates a contribution to the non-collinearity between the photon and the leading jet;
- $\Delta\eta$ is the difference in pseudorapidity between the jet and the outgoing photon: $\Delta\eta = \eta^{\text{jet}} - \eta^\gamma$, where η^{jet} and η^γ denote the pseudorapidity of the jet and the photon, respectively. This variable is sensitive to the dynamical properties of the scattering process;
- $\Delta\phi^{e,\gamma}$ is the azimuthal angle between the scattered electron and the outgoing photon: $\Delta\phi^{e,\gamma} = |\phi^e - \phi^\gamma|$, where ϕ^e denotes the azimuthal angle of the electron; this and the following variable are sensitive to higher-order processes and to whether the process is LL or QQ;
- $\Delta\eta^{e,\gamma}$ is the difference in pseudorapidity between the scattered electron and the photon: $\Delta\eta^{e,\gamma} = \eta^e - \eta^\gamma$, where η^e denotes the pseudorapidity of the electron.

A similar ZEUS analysis has been previously performed for photoproduction [44], studying all the present variables except those associated with the scattered electron.

5 Event simulation

Monte Carlo (MC) event samples were generated to evaluate the detector acceptance and to provide signal and background distributions. The program PYTHIA 6.416 [20] was used to simulate prompt-photon emission for the study of the event-reconstruction efficiency. In PYTHIA, this process is simulated as a DIS process with additional photon radiation from the quark line to account for QQ photons. Radiation from the lepton is not simulated.

The LL photons that were radiated into the detector and were isolated from the outgoing electron were simulated using the generator DJANGO 6 [47], an interface to the MC program HERACLES 4.6.6 [48]; higher-order QCD effects were included using the colour dipole model of ARIADNE 4.12 [49]. Hadronisation of the partonic final state was in each case performed by JETSET 7.4 [50] using the Lund string model [51]. Interference between the LL and QQ terms was neglected.

The main background to the QQ and LL photons came from photonic decays of neutral mesons produced in general DIS processes. This background was simulated using DJANGO 6, within the same framework as the LL events. This provided a realistic spectrum of single and multiple mesons with well modelled kinematic distributions.

The generated MC events were passed through ZEUS detector and trigger simulation programs based on GEANT 3.21 [52]. They were then reconstructed and analysed by the same programs as the data.

6 Theoretical calculations

The PYTHIA predictions and the predictions of two parton-level models were compared to the results of the present analysis. The NLO QCD calculation of Aurenche, Fontannaz and Guillet (AFG) [21], was performed in the $\overline{\text{MS}}$ scheme. Uncertainties on the QCD scale at this order contribute a normalisation uncertainty of typically $\pm 8\%$. This calculation was performed in the centre-of-mass frame and transformed into the laboratory frame, which introduces uncertainties on the cross sections in some regions of the parameter space due to non-perturbative effects [22]. The AFG predictions were calculated with a cut of 2.5 GeV on the photon transverse momentum in the centre-of-mass frame, and do not include an LL contribution, which was evaluated using the DJANGO-HERACLES simulation and added separately to the AFG calculation for comparison with the data. The uncertainties on the AFG predictions shown in the present paper represent the QCD scale uncertainties.

A calculation by Baranov, Lipatov and Zotov (BLZ) [23] used updated parameters for the present paper. It is based on the k_T -factorisation method. This approach uses unintegrated parton densities and takes into account both QQ and LL photons, neglecting the small interference contribution. The final result is obtained as the convolution of the off-shell scattering matrix element with the unintegrated quark distribution in the proton. In the k_T -factorisation theory, some part of the final-state jets can originate not only from the hard subprocess but also from the parton evolution cascade in the initial state. The quoted uncertainties on the BLZ predictions represent the QCD scale uncertainties.

In the previous ZEUS analysis of prompt photons in DIS, the measured variables were associated with the entire event, with the outgoing photon, and with jets. Comparisons were made to an earlier NLO QCD theory [53–55] and to BLZ. Both theories described the shapes of the single-particle cross sections well, but failed to reproduce the normalisation of the data. A later version of the original AFG calculation agreed well with the results [56], and has been used in the present study.

The predictions of AFG and BLZ were calculated at the parton level and incorporated kinematic and isolation criteria corresponding to the data. Corrections to the hadron level were made using PYTHIA to determine the ratio of the hadron-level cross sections to those at the parton level for each variable in each bin. The PYTHIA events were weighted at the parton level to represent the shapes of the AFG and BLZ distributions in x_γ^{meas} in order to calculate the hadronisation corrections for all the other measured variables. The corrections for AFG and BLZ were similar to within 10%. This procedure was also applied separately to the AFG predictions for the different Q^2 ranges.

For the BLZ x_γ^{meas} distribution, 98% of the parton-level cross section is in the (0.9, 1.0) bin; consequently, for this variable a transfer matrix from the parton to the hadron level was calculated using PYTHIA. The same procedure was used for the AFG x_γ^{meas} distribution. The relevant transfer matrices for the other variables gave similar results to the reweighting procedure.

7 Extraction of the photon signal

The event sample selected according to the criteria described in section 3 was dominated by background from neutral meson decays; thus the photon signal was extracted statistically following the approach used in previous ZEUS analyses [11–13, 16, 17].

The photon signal was evaluated making use of the width of the BEMC energy-cluster corresponding to the photon candidate. This was calculated as the variable

$$\langle \delta Z \rangle = \sum_i E_i |Z_i - Z_{\text{cluster}}| / \left(w_{\text{cell}} \sum_i E_i \right),$$

where Z_i is the Z position of the centre of the i -th cell, Z_{cluster} is the centroid of the EFO cluster, w_{cell} is the width of the cell in the Z direction, and E_i is the energy recorded in the cell. The sum runs over all BEMC cells in the EFO.

The distributions of $\langle \delta Z \rangle$ for the full data set and the fitted MC are shown in figure 2. The $\langle \delta Z \rangle$ distribution exhibits a double-peaked structure with the first peak at ≈ 0.1 , associated with the photon signal, and a second peak at ≈ 0.5 , dominated by the $\pi^0 \rightarrow \gamma\gamma$ background.

The contribution of isolated-photon events was determined for each bin in each measured variable by a χ^2 fit to the $\langle \delta Z \rangle$ distribution in the range $0.05 < \langle \delta Z \rangle < 0.8$, using the LL and QQ signal and background MC distributions as described in section 5. The mean value of $\chi^2/\text{n.d.f}$ was 1.2. Compared to the earlier ZEUS publication [19], improvements have been made in the modelling of the shapes of the $\langle \delta Z \rangle$ distributions of the QQ and LL contributions, using a comparison between the shapes associated with the scattered

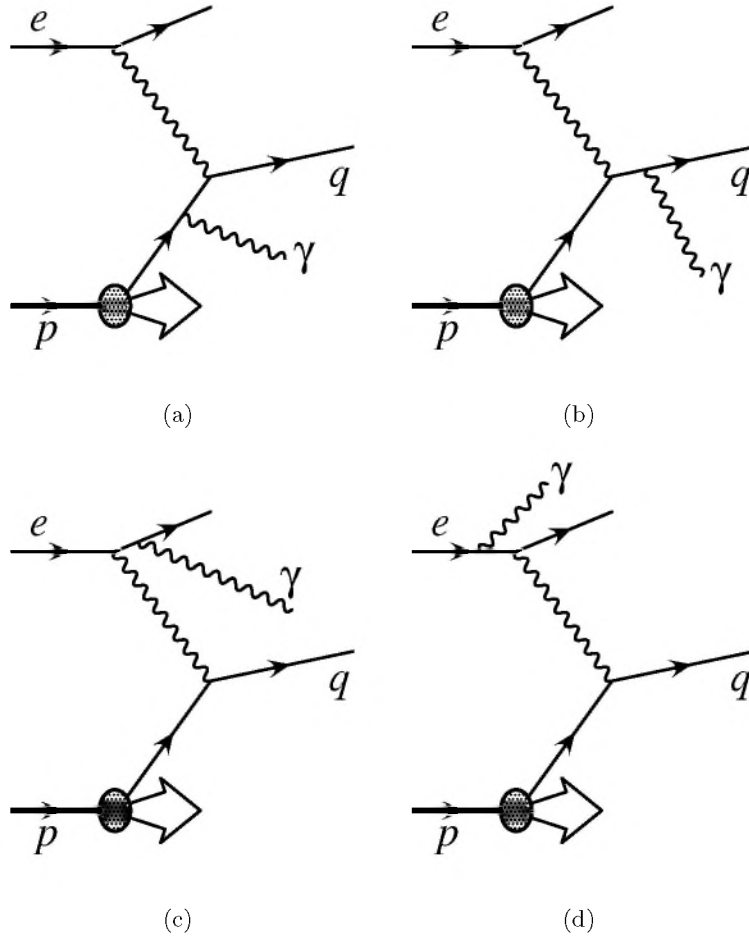


Figure 1. Lowest-order diagrams for photon production in ep scattering. (a), (b): quark radiative diagrams (QQ); (c), (d): lepton radiative diagrams (LL).

electron in MC simulation of DIS and in real data. By treating the LL and QQ photons separately, account is taken of the effect of their differing kinematic distributions on the acceptance, and the effect of their differing (η, E_T) distributions on the shape of the photon signal.

In performing the fit, the theoretically well determined LL contribution was kept constant at its MC-predicted value and the other components were varied. Of the 6149 events selected, 2451 ± 102 correspond to the extracted signal, including 526 LL photons. The fitted scale factor applied to the QQ contribution in figure 2 was 1.6, consistent with the earlier ZEUS analysis.

For a given observable Y , the production cross section was determined for each bin using

$$\frac{d\sigma}{dY} = \frac{\mathcal{A}_{QQ} \cdot N(\gamma_{QQ})}{\mathcal{L} \cdot \Delta Y} + \frac{d\sigma_{LL}^{MC}}{dY},$$

where $N(\gamma_{QQ})$ is the number of QQ photons extracted from the fit, ΔY is the bin width, \mathcal{L} is the total integrated luminosity, σ_{LL}^{MC} is the predicted cross section for LL photons from

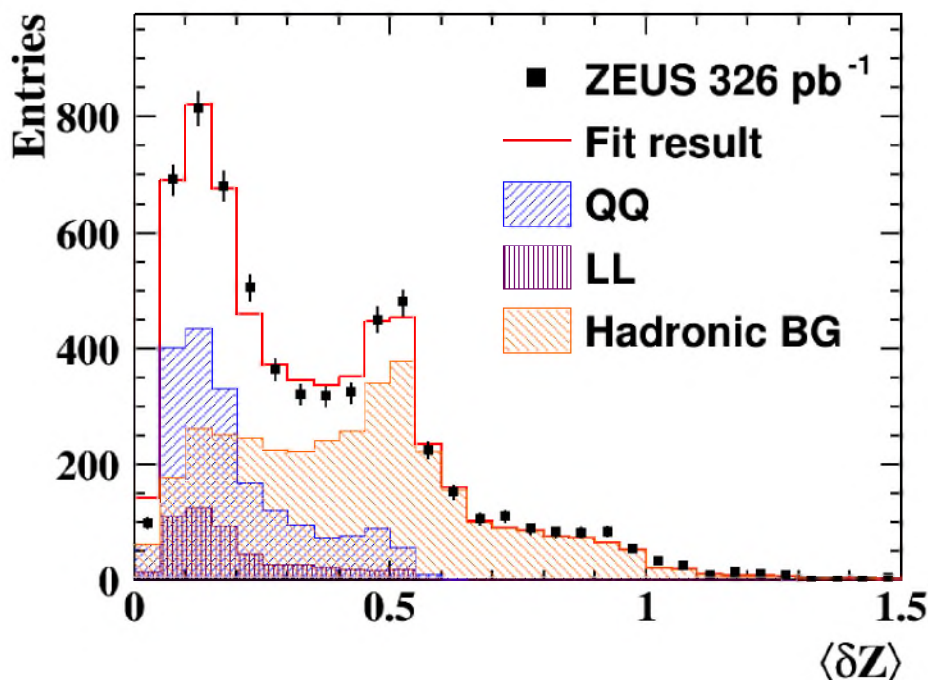


Figure 2. Distribution of $\langle \delta Z \rangle$ for the full data sample. The error bars represent the statistical uncertainties on the data points. The solid line shows a fit to the data of three components with fixed shapes as described in the text. The hatched histograms represent the LL and fitted QQ components of the fit and the fitted hadronic background (BG).

DJANGOHERACLES and \mathcal{A}_{QQ} is the acceptance correction for QQ photons. The value of \mathcal{A}_{QQ} was calculated, using the PYTHIA MC, from the ratio of the number of events generated to those reconstructed in a given bin; it lies in the range 0.91–2.28. To improve the representation of the data, and hence the accuracy of the acceptance corrections, the MC predictions were reweighted. This was done using parameterised functions of Q^2 and of η^γ , and also bin-by-bin as a function of photon energy; the three reweighting factors were applied multiplicatively. Their net effect on the acceptances was small.

8 Systematic uncertainties

The sources of systematic uncertainty on the measured cross sections are as in the previous paper [19]. The principal sources of uncertainty were evaluated as follows:

- the energy scale of the photon candidate was varied by $\pm 2\%$. The mean change of the cross section was $\pm 6\%$;
- the energy scale of the jets was varied by $\pm 1.5\%$ for jets with $E_T^{\text{jet}} > 10$ GeV, $\pm 2.5\%$ for jets with E_T^{jet} in the range $[6, 10]$ GeV and $\pm 4\%$ for jets with $E_T^{\text{jet}} < 6$ GeV. The uncertainty was typically $\pm 7\%$;
- the energy scale of the scattered electron was varied by $\pm 2\%$. The overall average effect on the cross sections was less than $\pm 1\%$.

Systematic uncertainties related to the MC generators were evaluated as follows:

- the dependence on the modelling of the hadronic background by means of DJANGO-HERACLES was investigated by varying the upper limit for the $\langle\delta Z\rangle$ fit in the range $[0.6, 1.0]$, giving variations that were typically $\pm 5\%$;
- uncertainties in the acceptance due to the PYTHIA model were accounted for by taking half of the change attributable to the reweighting described in section 7 as a systematic uncertainty; for most bins the effect was approximately 1%.

Other sources of systematic uncertainty were found to be negligible and were ignored [17, 57]: these included variations on the cuts on ΔR , the track momentum, $E - p_Z$, Z_{vtx} and the electromagnetic fraction of the photon shower, and a variation of 5% on the LL fraction.

The systematic uncertainties were symmetrised by taking the mean of the positive and negative uncertainty values and were combined in quadrature. The common uncertainty of 1.8% on the luminosity measurement is not included in the tables and figures.

9 Results

Differential cross sections for the production of an isolated photon in DIS with an additional jet have been measured in the laboratory frame in the kinematic region defined by $4 < E_T^\gamma < 15$ GeV, $-0.7 < \eta^\gamma < 0.9$, $E_T^{\text{jet}} > 2.5$ GeV and $-1.5 < \eta^{\text{jet}} < 1.8$. The DIS electron was constrained to be in the angular range $\theta_e > 140^\circ$, with energy greater than 10 GeV and $10 < Q^2 < 350$ GeV², where Q^2 was determined from the electron scattering angle. The jets were formed according to the k_T -clustering algorithm with the R parameter set to 1.0. Photon isolation was imposed such that at least 90% of the energy of the jet-like object containing the photon belonged to the photon.

The differential cross sections for the full Q^2 range as functions of x_γ^{meas} , x_p^{obs} , $\Delta\phi$, $\Delta\eta$, $\Delta\phi^{e,\gamma}$ and $\Delta\eta^{e,\gamma}$ are shown in figure 3 and are given in tables 1–6, which also list the values of the LL contributions and the hadronisation corrections. The cross section decreases with increasing x_p^{obs} , having a peak around 0.01, and rises at high values of x_γ^{meas} , $\Delta\phi$ and $\Delta\phi^{e,\gamma}$. The predictions for the sum of the expected LL contribution from DJANGO-HERACLES and a factor of 1.6 times the expected QQ contribution from PYTHIA agree well with the measurements. The success of the PYTHIA calculation can be attributed to its use of a leading-logarithm approach to gluon emission to augment its LO parton-scattering calculation.

The differential cross sections for the separate ranges $10 < Q^2 < 30$ GeV² and $30 < Q^2 < 350$ GeV² are shown in figures 4 and 5. In both these ranges, a good description of the data is given by the combination of the LL and PYTHIA MCs. The LL contribution is small in the lower Q^2 region, as was already seen in figure 3(a) of the earlier ZEUS publication [19]. In the higher Q^2 range, the LL component contributes significantly, as can be seen in the x_p^{obs} , $\Delta\phi$, $\Delta\eta$, and $\Delta\eta^{e,\gamma}$ distributions where it is dominant at high values of these variables. This reflects the changes with Q^2 in the structure of the contributing processes.

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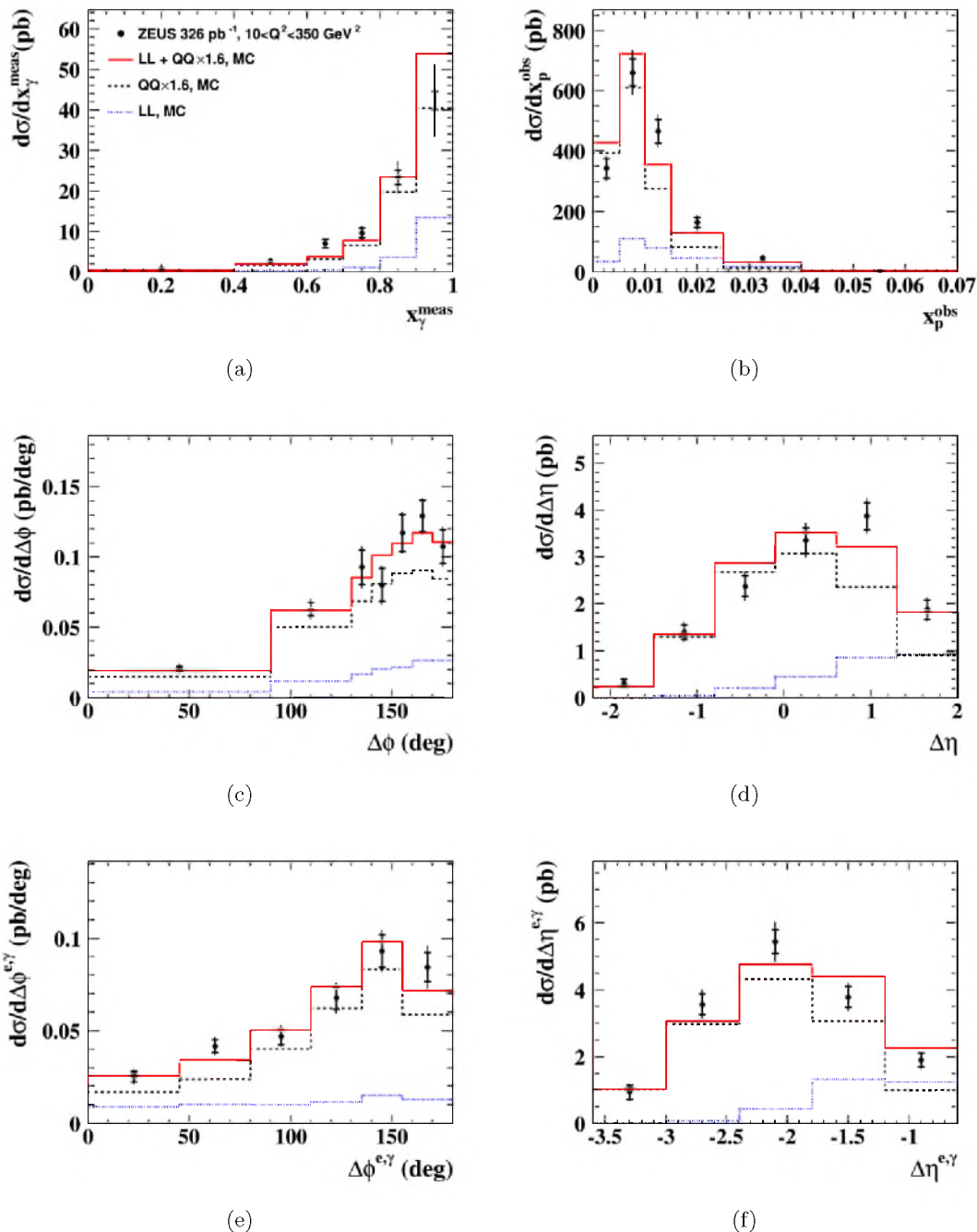


Figure 3. Differential cross sections in (a) x_γ^{meas} , (b) x_p^{obs} , (c) $\Delta\phi$, (d) $\Delta\eta$, (e) $\Delta\phi^{e,\gamma}$, and (f) $\Delta\eta^{e,\gamma}$, for the full range $10 < Q^2 < 350 \text{ GeV}^2$. The inner and outer error bars show, respectively, the statistical uncertainty and the statistical and systematic uncertainties added in quadrature. The solid histograms are the Monte Carlo predictions from the sum of QQ photons from PYTHIA normalised by a factor 1.6 plus DJANGO-HERACLES LL photons. The dashed (dotted) lines show the QQ (LL) contributions.

x_γ^{meas} range	$\frac{d\sigma}{dx_\gamma^{\text{meas}}} \text{ (pb)}$	$\frac{d\sigma_{\text{LL}}}{dx_\gamma^{\text{meas}}} \text{ (pb)}$	had. cor.
$10 < Q^2 < 350 \text{ GeV}^2$			
0.0 – 0.4	$0.94 \pm 0.20(\text{stat.}) \pm 0.11(\text{sys.})$	$0.06 \pm 0.01(\text{stat.})$	0.63
0.4 – 0.6	$2.73 \pm 0.43(\text{stat.}) \pm 0.32(\text{sys.})$	$0.29 \pm 0.04(\text{stat.})$	0.90
0.6 – 0.7	$7.06 \pm 1.14(\text{stat.}) \pm 0.38(\text{sys.})$	$0.65 \pm 0.09(\text{stat.})$	1.27
0.7 – 0.8	$9.64 \pm 1.24(\text{stat.}) \pm 1.06(\text{sys.})$	$1.17 \pm 0.12(\text{stat.})$	1.93
0.8 – 0.9	$23.40 \pm 1.75(\text{stat.}) \pm 3.51(\text{sys.})$	$3.67 \pm 0.22(\text{stat.})$	2.06
0.9 – 1.0	$42.34 \pm 2.26(\text{stat.}) \pm 8.54(\text{sys.})$	$13.49 \pm 0.42(\text{stat.})$	0.64
$10 < Q^2 < 30 \text{ GeV}^2$			
0.0 – 0.4	$0.45 \pm 0.15(\text{stat.}) \pm 0.09(\text{sys.})$	$0.01 \pm 0.01(\text{stat.})$	0.68
0.4 – 0.6	$1.19 \pm 0.31(\text{stat.}) \pm 0.18(\text{sys.})$	$0.07 \pm 0.02(\text{stat.})$	1.00
0.6 – 0.7	$4.30 \pm 0.88(\text{stat.}) \pm 0.49(\text{sys.})$	$0.23 \pm 0.06(\text{stat.})$	1.30
0.7 – 0.8	$5.58 \pm 0.88(\text{stat.}) \pm 0.69(\text{sys.})$	$0.16 \pm 0.04(\text{stat.})$	2.02
0.8 – 0.9	$9.27 \pm 1.20(\text{stat.}) \pm 1.32(\text{sys.})$	$0.54 \pm 0.08(\text{stat.})$	2.11
0.9 – 1.0	$17.76 \pm 1.37(\text{stat.}) \pm 3.73(\text{sys.})$	$1.89 \pm 0.16(\text{stat.})$	0.63
$30 \leq Q^2 < 350 \text{ GeV}^2$			
0.0 – 0.4	$0.38 \pm 0.15(\text{stat.}) \pm 0.05(\text{sys.})$	$0.06 \pm 0.01(\text{stat.})$	0.60
0.4 – 0.6	$1.55 \pm 0.30(\text{stat.}) \pm 0.23(\text{sys.})$	$0.22 \pm 0.04(\text{stat.})$	0.82
0.6 – 0.7	$2.50 \pm 0.73(\text{stat.}) \pm 0.36(\text{sys.})$	$0.42 \pm 0.07(\text{stat.})$	1.25
0.7 – 0.8	$4.15 \pm 0.89(\text{stat.}) \pm 0.53(\text{sys.})$	$1.01 \pm 0.11(\text{stat.})$	1.86
0.8 – 0.9	$13.90 \pm 1.27(\text{stat.}) \pm 2.01(\text{sys.})$	$3.14 \pm 0.20(\text{stat.})$	2.02
0.9 – 1.0	$25.81 \pm 1.89(\text{stat.}) \pm 4.74(\text{sys.})$	$11.61 \pm 0.38(\text{stat.})$	0.65

Table 1. Measured differential cross-section $\frac{d\sigma}{dx_\gamma^{\text{meas}}}$. The quoted systematic uncertainty includes all the components added in quadrature. The calculated LL contribution which was added to the PYTHIA and AFG calculations is also listed, and the hadronisation correction calculated for the AFG predictions. Differences between cross sections in the first section and the sum of the corresponding values in the second and third sections are of statistical origin.

x_p^{obs} range	$\frac{d\sigma}{dx_p^{\text{obs}}} \text{ (pb)}$	$\frac{d\sigma_{\text{LL}}}{dx_p^{\text{obs}}} \text{ (pb)}$	had. cor.
$10 < Q^2 < 350 \text{ GeV}^2$			
0.000 – 0.005	$344.3 \pm 31.7(\text{stat.}) \pm 22.9(\text{sys.})$	$35.2 \pm 3.0(\text{stat.})$	0.69
0.005 – 0.010	$661.8 \pm 45.3(\text{stat.}) \pm 56.6(\text{sys.})$	$110.8 \pm 5.3(\text{stat.})$	0.81
0.010 – 0.015	$467.1 \pm 38.9(\text{stat.}) \pm 35.5(\text{sys.})$	$80.0 \pm 4.5(\text{stat.})$	0.91
0.015 – 0.025	$164.5 \pm 16.5(\text{stat.}) \pm 16.1(\text{sys.})$	$46.6 \pm 2.4(\text{stat.})$	0.99
0.025 – 0.040	$46.7 \pm 6.8 (\text{stat.}) \pm 2.7 (\text{sys.})$	$18.7 \pm 1.3(\text{stat.})$	1.06
0.040 – 0.070	$3.3 \pm 0.6 (\text{stat.}) \pm 2.1 (\text{sys.})$	$3.3 \pm 0.4(\text{stat.})$	1.00
$10 < Q^2 < 30 \text{ GeV}^2$			
0.000 – 0.005	$201.8 \pm 25.0(\text{stat.}) \pm 11.1(\text{sys.})$	$8.4 \pm 1.4(\text{stat.})$	0.71
0.005 – 0.010	$319.6 \pm 31.4(\text{stat.}) \pm 31.8(\text{sys.})$	$19.4 \pm 2.2(\text{stat.})$	0.84
0.010 – 0.015	$195.5 \pm 24.5(\text{stat.}) \pm 20.5(\text{sys.})$	$12.7 \pm 1.8(\text{stat.})$	0.98
0.015 – 0.025	$68.1 \pm 10.4(\text{stat.}) \pm 9.8 (\text{sys.})$	$5.6 \pm 0.9(\text{stat.})$	1.03
0.025 – 0.040	$18.7 \pm 4.1 (\text{stat.}) \pm 9.5 (\text{sys.})$	$2.1 \pm 0.4(\text{stat.})$	1.08
0.040 – 0.070	$0.2 \pm 0.1 (\text{stat.}) \pm 0.1 (\text{sys.})$	$0.2 \pm 0.1(\text{stat.})$	0.95
$30 \leq Q^2 < 350 \text{ GeV}^2$			
0.000 – 0.005	$149.3 \pm 20.0(\text{stat.}) \pm 9.1 (\text{sys.})$	$26.8 \pm 2.6(\text{stat.})$	0.68
0.005 – 0.010	$340.7 \pm 32.9(\text{stat.}) \pm 25.0(\text{sys.})$	$91.4 \pm 4.8(\text{stat.})$	0.78
0.010 – 0.015	$271.7 \pm 30.5(\text{stat.}) \pm 17.4(\text{sys.})$	$67.3 \pm 4.1(\text{stat.})$	0.88
0.015 – 0.025	$97.7 \pm 12.8(\text{stat.}) \pm 8.1 (\text{sys.})$	$41.0 \pm 2.3(\text{stat.})$	0.97
0.025 – 0.040	$37.5 \pm 5.3 (\text{stat.}) \pm 3.1 (\text{sys.})$	$16.6 \pm 1.2(\text{stat.})$	1.06
0.040 – 0.070	$3.0 \pm 1.0 (\text{stat.}) \pm 2.1 (\text{sys.})$	$3.0 \pm 0.4(\text{stat.})$	1.01

Table 2. Measured differential cross-section $\frac{d\sigma}{dx_p^{\text{obs}}}$. Details as in table 1.

$\Delta\phi$ range (deg)	$\frac{d\sigma}{d\Delta\phi}$ (pb/deg)	$\frac{d\sigma_{\text{LL}}}{d\Delta\phi}$ (pb/deg)	had. cor.
$10 < Q^2 < 350 \text{ GeV}^2$			
0 – 90	$0.020 \pm 0.002(\text{stat.}) \pm 0.003(\text{sys.})$	$0.004 \pm 0.001(\text{stat.})$	0.68
90 – 130	$0.063 \pm 0.005(\text{stat.}) \pm 0.005(\text{sys.})$	$0.012 \pm 0.001(\text{stat.})$	0.82
130 – 140	$0.093 \pm 0.012(\text{stat.}) \pm 0.008(\text{sys.})$	$0.017 \pm 0.002(\text{stat.})$	0.88
140 – 150	$0.080 \pm 0.012(\text{stat.}) \pm 0.007(\text{sys.})$	$0.021 \pm 0.002(\text{stat.})$	0.92
150 – 160	$0.117 \pm 0.013(\text{stat.}) \pm 0.006(\text{sys.})$	$0.021 \pm 0.002(\text{stat.})$	0.95
160 – 170	$0.129 \pm 0.011(\text{stat.}) \pm 0.005(\text{sys.})$	$0.027 \pm 0.002(\text{stat.})$	0.95
170 – 180	$0.108 \pm 0.012(\text{stat.}) \pm 0.007(\text{sys.})$	$0.026 \pm 0.002(\text{stat.})$	0.94
$10 < Q^2 < 30 \text{ GeV}^2$			
0 – 90	$0.004 \pm 0.001(\text{stat.}) \pm 0.001(\text{sys.})$	$0.000 \pm 0.001(\text{stat.})$	0.68
90 – 130	$0.023 \pm 0.003(\text{stat.}) \pm 0.002(\text{sys.})$	$0.001 \pm 0.001(\text{stat.})$	0.78
130 – 140	$0.042 \pm 0.010(\text{stat.}) \pm 0.007(\text{sys.})$	$0.003 \pm 0.001(\text{stat.})$	0.79
140 – 150	$0.047 \pm 0.009(\text{stat.}) \pm 0.005(\text{sys.})$	$0.004 \pm 0.001(\text{stat.})$	0.85
150 – 160	$0.057 \pm 0.010(\text{stat.}) \pm 0.003(\text{sys.})$	$0.005 \pm 0.001(\text{stat.})$	0.91
160 – 170	$0.079 \pm 0.009(\text{stat.}) \pm 0.004(\text{sys.})$	$0.007 \pm 0.001(\text{stat.})$	0.93
170 – 180	$0.064 \pm 0.009(\text{stat.}) \pm 0.005(\text{sys.})$	$0.007 \pm 0.001(\text{stat.})$	0.93
$30 \leq Q^2 < 350 \text{ GeV}^2$			
0 – 90	$0.015 \pm 0.002(\text{stat.}) \pm 0.002(\text{sys.})$	$0.004 \pm 0.001(\text{stat.})$	0.68
90 – 130	$0.040 \pm 0.004(\text{stat.}) \pm 0.003(\text{sys.})$	$0.011 \pm 0.001(\text{stat.})$	0.83
130 – 140	$0.049 \pm 0.008(\text{stat.}) \pm 0.002(\text{sys.})$	$0.014 \pm 0.001(\text{stat.})$	0.96
140 – 150	$0.030 \pm 0.008(\text{stat.}) \pm 0.001(\text{sys.})$	$0.017 \pm 0.002(\text{stat.})$	0.99
150 – 160	$0.064 \pm 0.009(\text{stat.}) \pm 0.007(\text{sys.})$	$0.016 \pm 0.001(\text{stat.})$	1.01
160 – 170	$0.046 \pm 0.007(\text{stat.}) \pm 0.005(\text{sys.})$	$0.020 \pm 0.002(\text{stat.})$	1.01
170 – 180	$0.045 \pm 0.009(\text{stat.}) \pm 0.003(\text{sys.})$	$0.019 \pm 0.002(\text{stat.})$	0.97

Table 3. Measured differential cross-section $\frac{d\sigma}{d\Delta\phi}$. Details as in table 1.

$\Delta\eta$ range	$\frac{d\sigma}{d\Delta\eta}$ (pb)	$\frac{d\sigma_{LL}}{d\Delta\eta}$ (pb)	had. cor.
$10 < Q^2 < 350 \text{ GeV}^2$			
$-2.2 - -1.5$	$0.32 \pm 0.08(\text{stat.}) \pm 0.05(\text{sys.})$	$0.01 \pm 0.01(\text{stat.})$	0.76
$-1.5 - -0.8$	$1.41 \pm 0.15(\text{stat.}) \pm 0.14(\text{sys.})$	$0.06 \pm 0.01(\text{stat.})$	0.66
$-0.8 - -0.1$	$2.38 \pm 0.22(\text{stat.}) \pm 0.21(\text{sys.})$	$0.21 \pm 0.02(\text{stat.})$	0.74
$-0.1 - 0.6$	$3.36 \pm 0.27(\text{stat.}) \pm 0.23(\text{sys.})$	$0.45 \pm 0.03(\text{stat.})$	0.87
$0.6 - 1.3$	$3.88 \pm 0.28(\text{stat.}) \pm 0.22(\text{sys.})$	$0.87 \pm 0.04(\text{stat.})$	1.04
$1.3 - 2.0$	$1.88 \pm 0.21(\text{stat.}) \pm 0.12(\text{sys.})$	$0.92 \pm 0.04(\text{stat.})$	1.11
$10 < Q^2 < 30 \text{ GeV}^2$			
$-2.2 - -1.5$	$0.14 \pm 0.05(\text{stat.}) \pm 0.03(\text{sys.})$	$0.00 \pm 0.01(\text{stat.})$	0.63
$-1.5 - -0.8$	$0.51 \pm 0.12(\text{stat.}) \pm 0.04(\text{sys.})$	$0.00 \pm 0.01(\text{stat.})$	0.68
$-0.8 - -0.1$	$1.16 \pm 0.15(\text{stat.}) \pm 0.09(\text{sys.})$	$0.04 \pm 0.01(\text{stat.})$	0.77
$-0.1 - 0.6$	$1.70 \pm 0.19(\text{stat.}) \pm 0.15(\text{sys.})$	$0.08 \pm 0.01(\text{stat.})$	0.90
$0.6 - 1.3$	$1.67 \pm 0.19(\text{stat.}) \pm 0.13(\text{sys.})$	$0.14 \pm 0.02(\text{stat.})$	1.08
$1.3 - 2.0$	$0.71 \pm 0.13(\text{stat.}) \pm 0.07(\text{sys.})$	$0.13 \pm 0.02(\text{stat.})$	1.07
$30 \leq Q^2 < 350 \text{ GeV}^2$			
$-2.2 - -1.5$	$0.20 \pm 0.07(\text{stat.}) \pm 0.03(\text{sys.})$	$0.00 \pm 0.01(\text{stat.})$	0.83
$-1.5 - -0.8$	$0.86 \pm 0.09(\text{stat.}) \pm 0.09(\text{sys.})$	$0.05 \pm 0.01(\text{stat.})$	0.65
$-0.8 - -0.1$	$1.25 \pm 0.16(\text{stat.}) \pm 0.13(\text{sys.})$	$0.16 \pm 0.02(\text{stat.})$	0.72
$-0.1 - 0.6$	$1.68 \pm 0.19(\text{stat.}) \pm 0.08(\text{sys.})$	$0.37 \pm 0.03(\text{stat.})$	0.85
$0.6 - 1.3$	$2.23 \pm 0.22(\text{stat.}) \pm 0.19(\text{sys.})$	$0.72 \pm 0.04(\text{stat.})$	1.02
$1.3 - 2.0$	$1.16 \pm 0.16(\text{stat.}) \pm 0.06(\text{sys.})$	$0.80 \pm 0.04(\text{stat.})$	1.14

Table 4. Measured differential cross-section $\frac{d\sigma}{d\Delta\eta}$. Details as in table 1.

$\Delta\phi^{e,\gamma}$ range (deg)	$\frac{d\sigma}{d\Delta\phi^{e,\gamma}}$ (pb/deg)	$\frac{d\sigma_{\text{LL}}}{d\Delta\phi^{e,\gamma}}$ (pb/deg)	had. cor.
$10 < Q^2 < 350 \text{ GeV}^2$			
0 – 45	$0.025 \pm 0.003(\text{stat.}) \pm 0.002(\text{sys.})$	$0.009 \pm 0.001(\text{stat.})$	0.95
45 – 80	$0.042 \pm 0.004(\text{stat.}) \pm 0.003(\text{sys.})$	$0.010 \pm 0.001(\text{stat.})$	0.94
80 – 110	$0.047 \pm 0.004(\text{stat.}) \pm 0.003(\text{sys.})$	$0.010 \pm 0.001(\text{stat.})$	0.92
110 – 135	$0.068 \pm 0.006(\text{stat.}) \pm 0.006(\text{sys.})$	$0.012 \pm 0.001(\text{stat.})$	0.85
135 – 155	$0.093 \pm 0.009(\text{stat.}) \pm 0.007(\text{sys.})$	$0.015 \pm 0.001(\text{stat.})$	0.79
155 – 180	$0.085 \pm 0.008(\text{stat.}) \pm 0.008(\text{sys.})$	$0.013 \pm 0.001(\text{stat.})$	0.73
$10 < Q^2 < 30 \text{ GeV}^2$			
0 – 45	$0.013 \pm 0.002(\text{stat.}) \pm 0.002(\text{sys.})$	$0.002 \pm 0.001(\text{stat.})$	0.95
45 – 80	$0.018 \pm 0.003(\text{stat.}) \pm 0.001(\text{sys.})$	$0.002 \pm 0.001(\text{stat.})$	0.94
80 – 110	$0.024 \pm 0.003(\text{stat.}) \pm 0.002(\text{sys.})$	$0.001 \pm 0.001(\text{stat.})$	0.91
110 – 135	$0.033 \pm 0.005(\text{stat.}) \pm 0.002(\text{sys.})$	$0.002 \pm 0.001(\text{stat.})$	0.85
135 – 155	$0.031 \pm 0.006(\text{stat.}) \pm 0.002(\text{sys.})$	$0.001 \pm 0.001(\text{stat.})$	0.78
155 – 180	$0.038 \pm 0.005(\text{stat.}) \pm 0.004(\text{sys.})$	$0.002 \pm 0.001(\text{stat.})$	0.80
$30 \leq Q^2 < 350 \text{ GeV}^2$			
0 – 45	$0.012 \pm 0.002(\text{stat.}) \pm 0.001(\text{sys.})$	$0.007 \pm 0.001(\text{stat.})$	0.95
45 – 80	$0.024 \pm 0.002(\text{stat.}) \pm 0.002(\text{sys.})$	$0.009 \pm 0.001(\text{stat.})$	0.95
80 – 110	$0.023 \pm 0.003(\text{stat.}) \pm 0.002(\text{sys.})$	$0.009 \pm 0.001(\text{stat.})$	0.93
110 – 135	$0.036 \pm 0.004(\text{stat.}) \pm 0.003(\text{sys.})$	$0.010 \pm 0.001(\text{stat.})$	0.86
135 – 155	$0.063 \pm 0.007(\text{stat.}) \pm 0.005(\text{sys.})$	$0.014 \pm 0.001(\text{stat.})$	0.80
155 – 180	$0.047 \pm 0.006(\text{stat.}) \pm 0.004(\text{sys.})$	$0.011 \pm 0.001(\text{stat.})$	0.70

Table 5. Measured differential cross-section $\frac{d\sigma}{d\Delta\phi^{e,\gamma}}$. Details as in table 1.

$\Delta\eta^{e,\gamma}$ range	$\frac{d\sigma}{d\Delta\eta^{e,\gamma}}(\text{pb})$	$\frac{d\sigma_{\text{I,L}}}{d\Delta\eta^{e,\gamma}}(\text{pb})$	had. cor.
$10 < Q^2 < 350 \text{ GeV}^2$			
$-3.6 - -3.0$	$0.94 \pm 0.21(\text{stat.}) \pm 0.12(\text{sys.})$	$0.02 \pm 0.01(\text{stat.})$	0.80
$-3.0 - -2.4$	$3.57 \pm 0.30(\text{stat.}) \pm 0.30(\text{sys.})$	$0.08 \pm 0.01(\text{stat.})$	0.82
$-2.4 - -1.8$	$5.44 \pm 0.36(\text{stat.}) \pm 0.45(\text{sys.})$	$0.45 \pm 0.03(\text{stat.})$	0.83
$-1.8 - -1.2$	$3.79 \pm 0.31(\text{stat.}) \pm 0.26(\text{sys.})$	$1.33 \pm 0.05(\text{stat.})$	0.85
$-1.2 - -0.6$	$1.90 \pm 0.21(\text{stat.}) \pm 0.11(\text{sys.})$	$1.24 \pm 0.05(\text{stat.})$	0.89
$10 < Q^2 < 30 \text{ GeV}^2$			
$-3.6 - -3.0$	$0.93 \pm 0.21(\text{stat.}) \pm 0.12(\text{sys.})$	$0.02 \pm 0.01(\text{stat.})$	0.81
$-3.0 - -2.4$	$2.60 \pm 0.25(\text{stat.}) \pm 0.19(\text{sys.})$	$0.06 \pm 0.01(\text{stat.})$	0.85
$-2.4 - -1.8$	$2.69 \pm 0.25(\text{stat.}) \pm 0.19(\text{sys.})$	$0.22 \pm 0.02(\text{stat.})$	0.89
$-1.8 - -1.2$	$0.86 \pm 0.15(\text{stat.}) \pm 0.07(\text{sys.})$	$0.19 \pm 0.02(\text{stat.})$	0.92
$30 \leq Q^2 < 350 \text{ GeV}^2$			
$-3.0 - -2.4$	$1.00 \pm 0.17(\text{stat.}) \pm 0.11(\text{sys.})$	$0.02 \pm 0.01(\text{stat.})$	0.77
$-2.4 - -1.8$	$2.72 \pm 0.26(\text{stat.}) \pm 0.25(\text{sys.})$	$0.23 \pm 0.02(\text{stat.})$	0.80
$-1.8 - -1.2$	$3.00 \pm 0.27(\text{stat.}) \pm 0.18(\text{sys.})$	$1.14 \pm 0.05(\text{stat.})$	0.84
$-1.2 - -0.6$	$1.90 \pm 0.21(\text{stat.}) \pm 0.11(\text{sys.})$	$1.24 \pm 0.05(\text{stat.})$	0.89

Table 6. Measured differential cross-section $\frac{d\sigma}{d\Delta\eta^{e,\gamma}}$. Details as in table 1.

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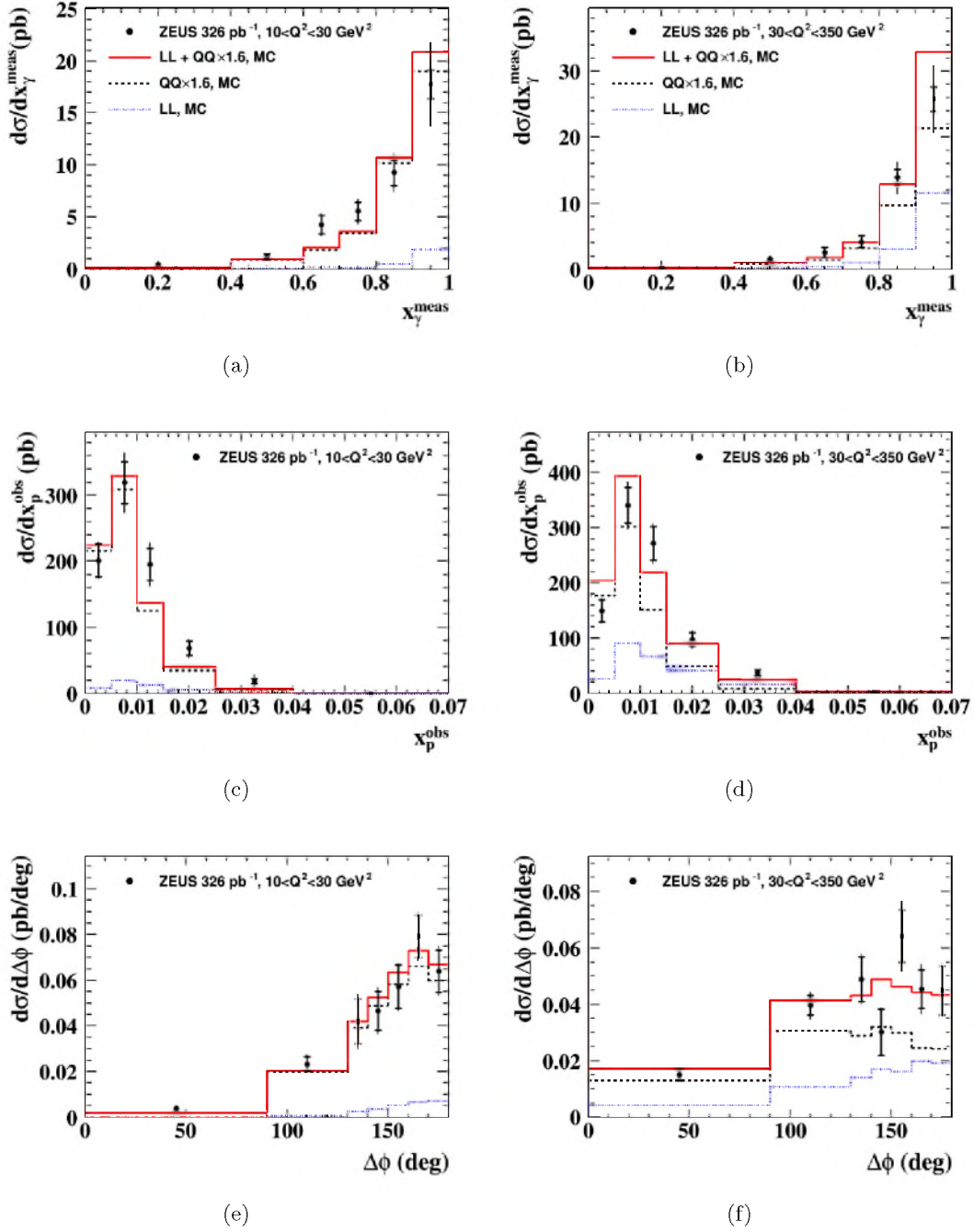


Figure 4. Differential cross sections for the regions $10 < Q^2 < 30$ and $30 < Q^2 < 350$ GeV^2 : (a, b) x_γ^{meas} , (c, d) x_p^{obs} , and (e, f) $\Delta\phi$. The inner and outer error bars show, respectively, the statistical uncertainty and the statistical and systematic uncertainties added in quadrature. The solid histograms are the Monte Carlo predictions from the sum of QQ photons from PYTHIA normalised by a factor 1.6 plus DJANGO-HERACLES LL photons. The dashed (dotted) lines show the QQ (LL) contributions.

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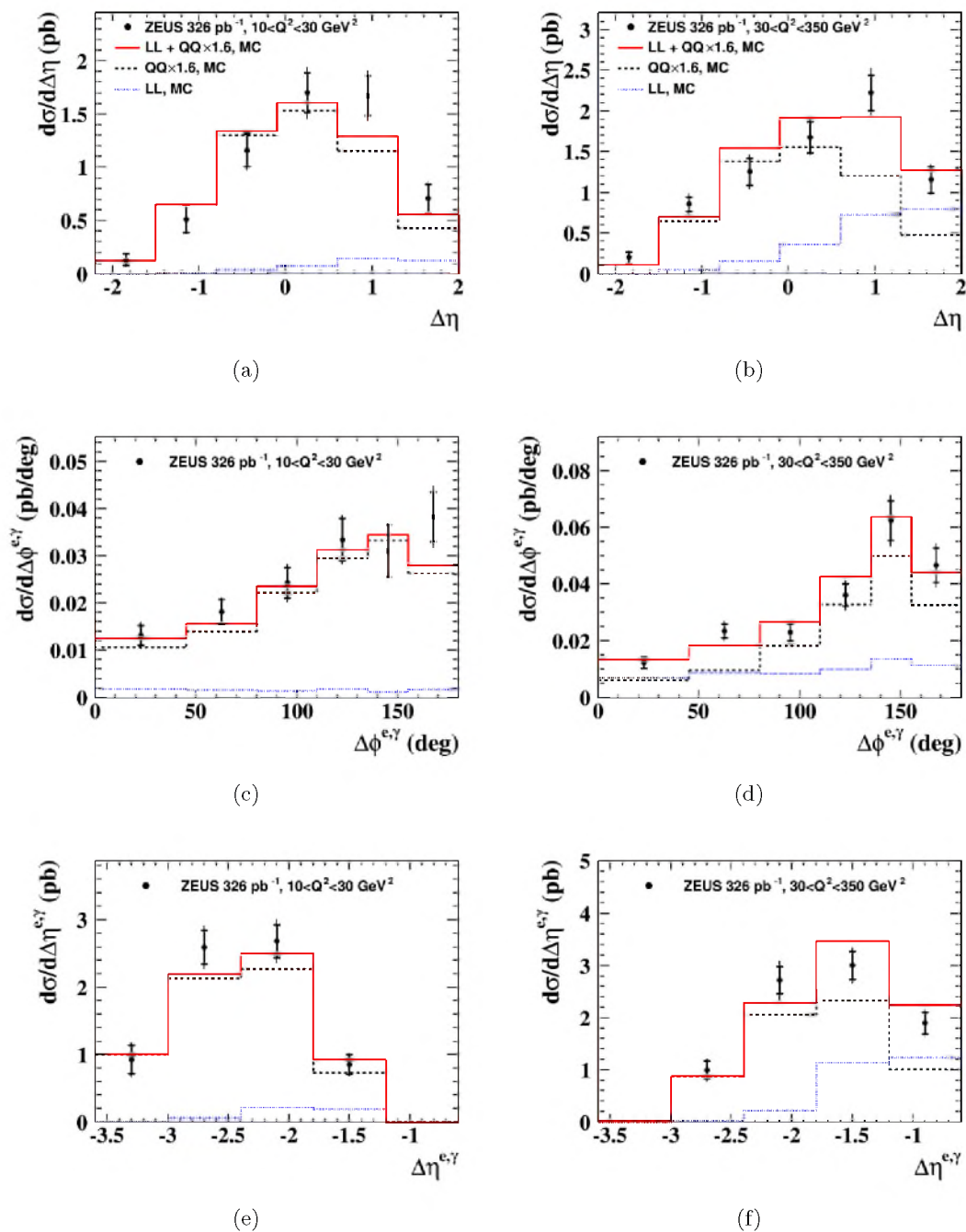


Figure 5. Differential cross sections for the regions $10 < Q^2 < 30$ and $30 < Q^2 < 350$ GeV^2 : (a, b) $\Delta\eta$, (c, d) $\Delta\phi^{e,\gamma}$, and (e, f) $\Delta\eta^{e,\gamma}$. The inner and outer error bars show, respectively, the statistical uncertainty and the statistical and systematic uncertainties added in quadrature. The solid histograms are the Monte Carlo predictions from the sum of QQ photons from PYTHIA normalised by a factor 1.6 plus DJANGOHERACLES LL photons. The dashed (dotted) lines show the QQ (LL) contributions.

The increased importance of the LL component at higher Q^2 is also reflected in the x_γ^{meas} distribution. Figure 6 presents the x_γ^{meas} and x_p^{obs} cross sections on a logarithmic scale. The data in the low- x_γ^{meas} region are satisfactorily described by PYTHIA without the need for further higher-order processes.

Comparisons of the data with the AFG and BLZ predictions are presented for the entire Q^2 range in figure 7. The updated BLZ predictions describe the shape of most of the distributions reasonably well, but there is an overestimation of about 20% in the overall cross section, and the extremely peaked prediction for the x_γ^{meas} distribution is not in agreement with the data. The AFG predictions describe all the distributions well and also agree in the overall normalisation.

Comparisons of the data with the AFG model in the two separate Q^2 ranges are shown in figures 8–9. In the higher Q^2 range, the description by AFG is excellent. In the lower range, the only deviation observable is in the $\Delta\eta$ distribution, where the data show a tendency towards higher values than the theory. This might be related to the cut of 2.5 GeV on the transverse photon momentum applied in the AFG calculation [21].

10 Summary

The production of isolated photons accompanied by jets has been measured in deep inelastic scattering with the ZEUS detector at HERA, using an integrated luminosity of 326 pb^{-1} . Expanding on earlier ZEUS results [19], which studied single-particle distributions, differential cross sections have been evaluated as functions of pairs of measured variables in combination. The kinematic region in the laboratory frame was defined by $4 < E_T^\gamma < 15 \text{ GeV}$, $-0.7 < \eta^\gamma < 0.9$, $E_T^{\text{jet}} > 2.5 \text{ GeV}$ and $-1.5 < \eta^{\text{jet}} < 1.8$. The DIS electron was constrained to be in the angular range $\theta_e > 140^\circ$, with energy greater than 10 GeV and $10 < Q^2 < 350 \text{ GeV}^2$, where Q^2 was determined from the electron scattering angle. The jets were formed according to the k_T -clustering algorithm with the R parameter set to 1.0. Photon isolation was imposed such that at least 90% of the energy of the jet-like object containing the photon belonged to the photon. Differential cross sections are presented for the following variables: the fraction of the incoming photon energy and momentum that is transferred to the outgoing photon and the leading jet; the fraction of the incoming proton energy transferred to the photon and leading jet; the differences in azimuthal angle and pseudorapidity between the outgoing photon and the leading jet and between the outgoing photon and the scattered electron.

The PYTHIA prediction for the quark-radiated photon component plus the DJANGO-HERACLES calculation for the lepton-radiated component describes all the distributions well if the PYTHIA prediction is scaled up by a factor of 1.6. This is also true if the data are divided into ranges above and below a value of $Q^2 = 30 \text{ GeV}^2$. Predictions from two theoretical models were also compared to the data. The BLZ model gives a fair description of the data but does not give a good description of the overall normalisation or the shape of some of the distributions. The AFG model gives an excellent description of the normalisation and almost all the distributions, both for the entire data set and for the separate Q^2 ranges.

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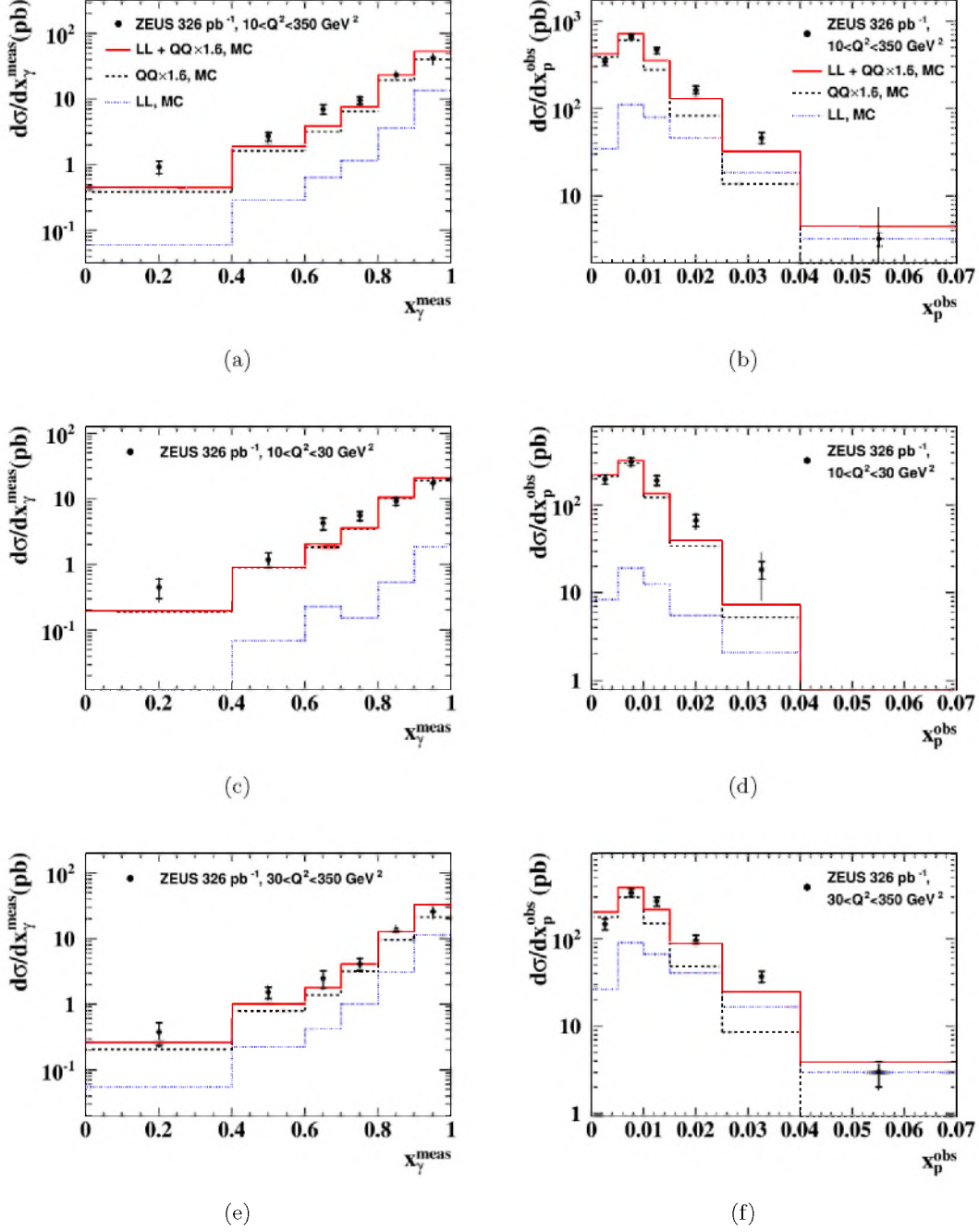


Figure 6. Differential cross sections in (a, c, e) x_γ^{meas} and (b, d, f) x_p^{obs} for (a, b) $10 < Q^2 < 350 \text{ GeV}^2$, (c, d) $10 < Q^2 < 30 \text{ GeV}^2$, and (e, f) $30 < Q^2 < 350 \text{ GeV}^2$. The distributions are as shown in figures 3–5 but with logarithmic vertical scale.

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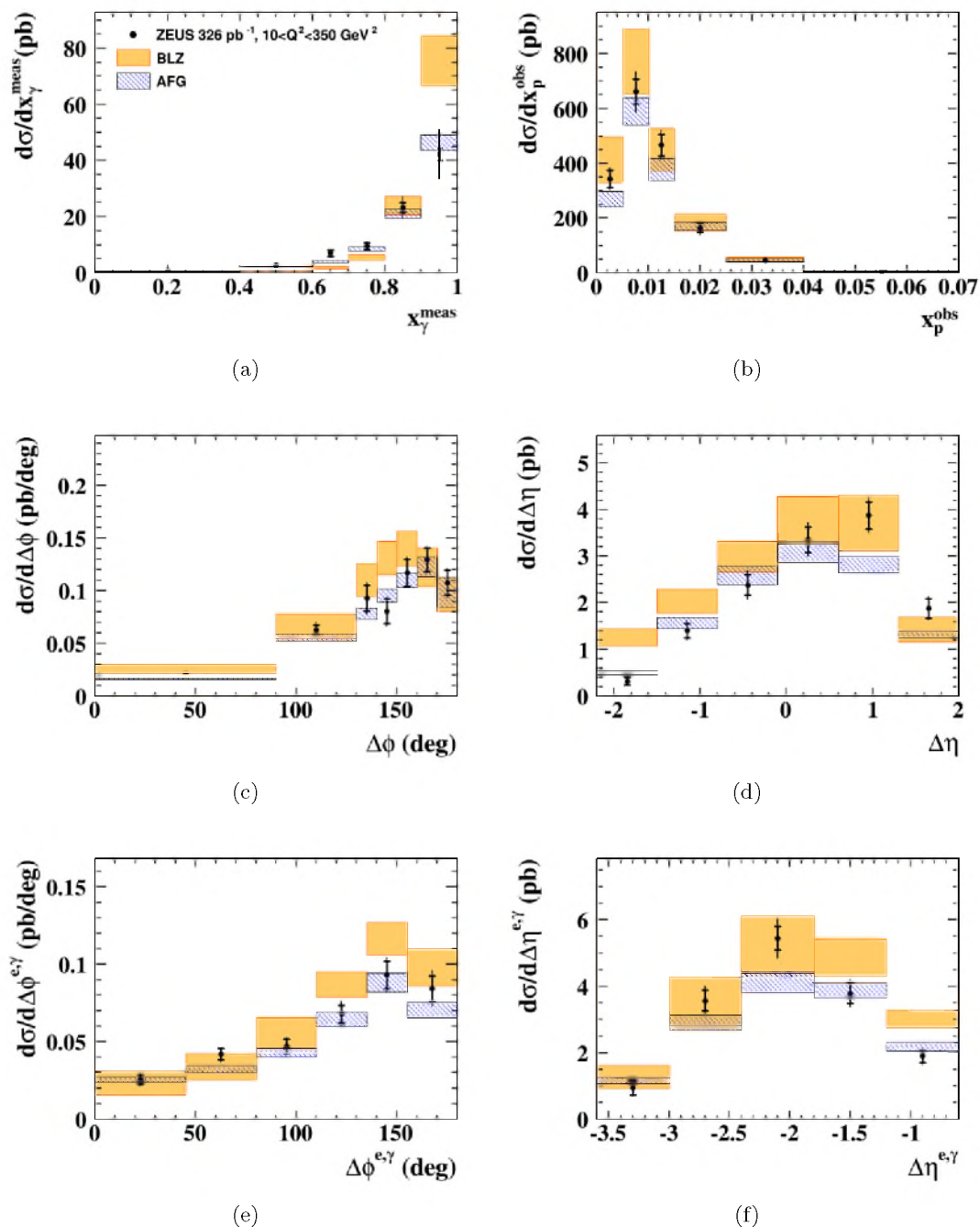


Figure 7. Differential cross sections for selected variables in the full Q^2 range $10 < Q^2 < 350 \text{ GeV}^2$: as in figure 3. Theoretical predictions from Aurenche et al. (AFG) and Baranov et al. (BLZ) are shown, with scale uncertainties indicated by the bands.

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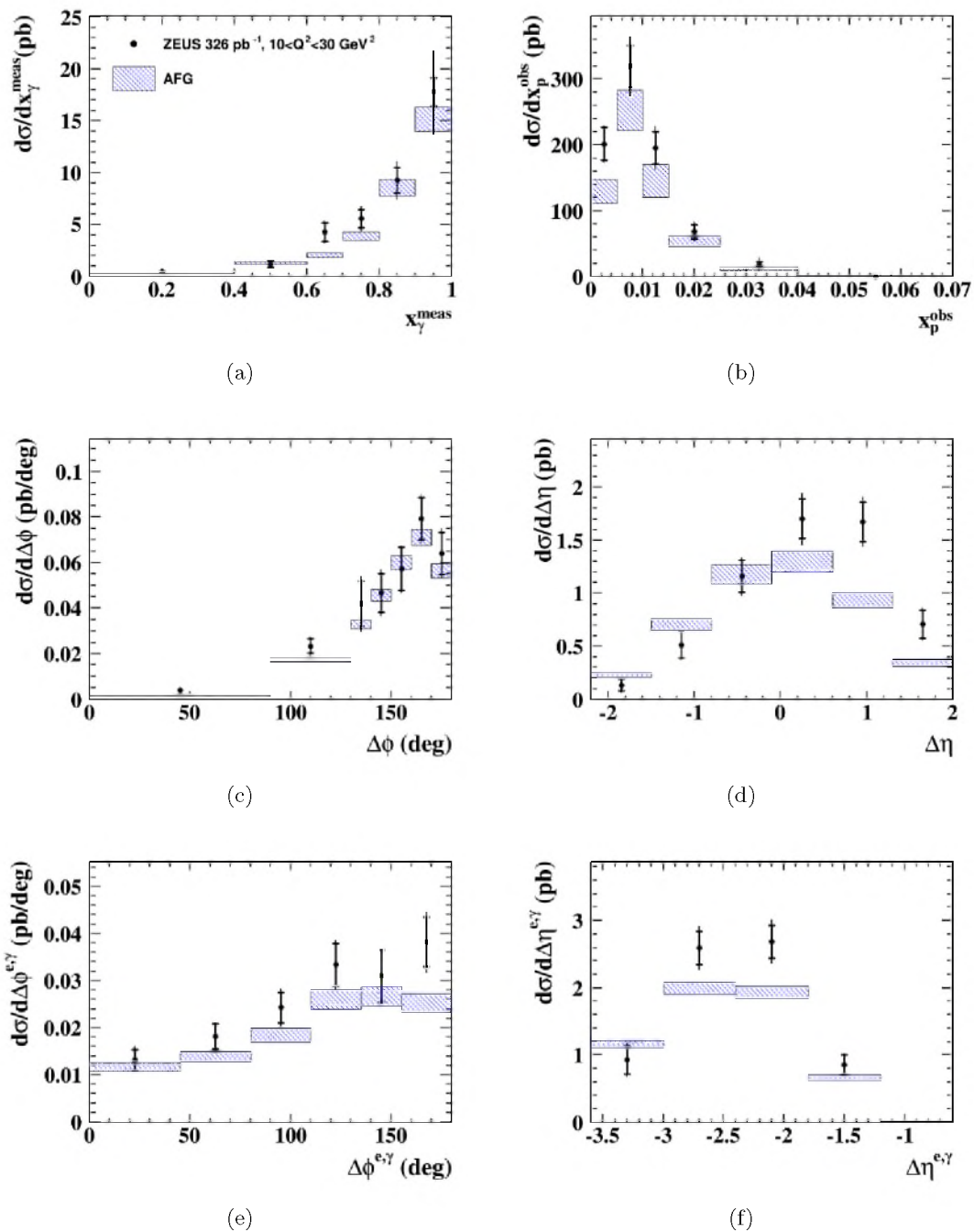


Figure 8. Differential cross sections for selected variables in the region $10 < Q^2 < 30 \text{ GeV}^2$ as in figures 4, 5. Theoretical predictions from Aurenche et al. (AFG) are shown, with associated uncertainties indicated by the bands.

ZEUS

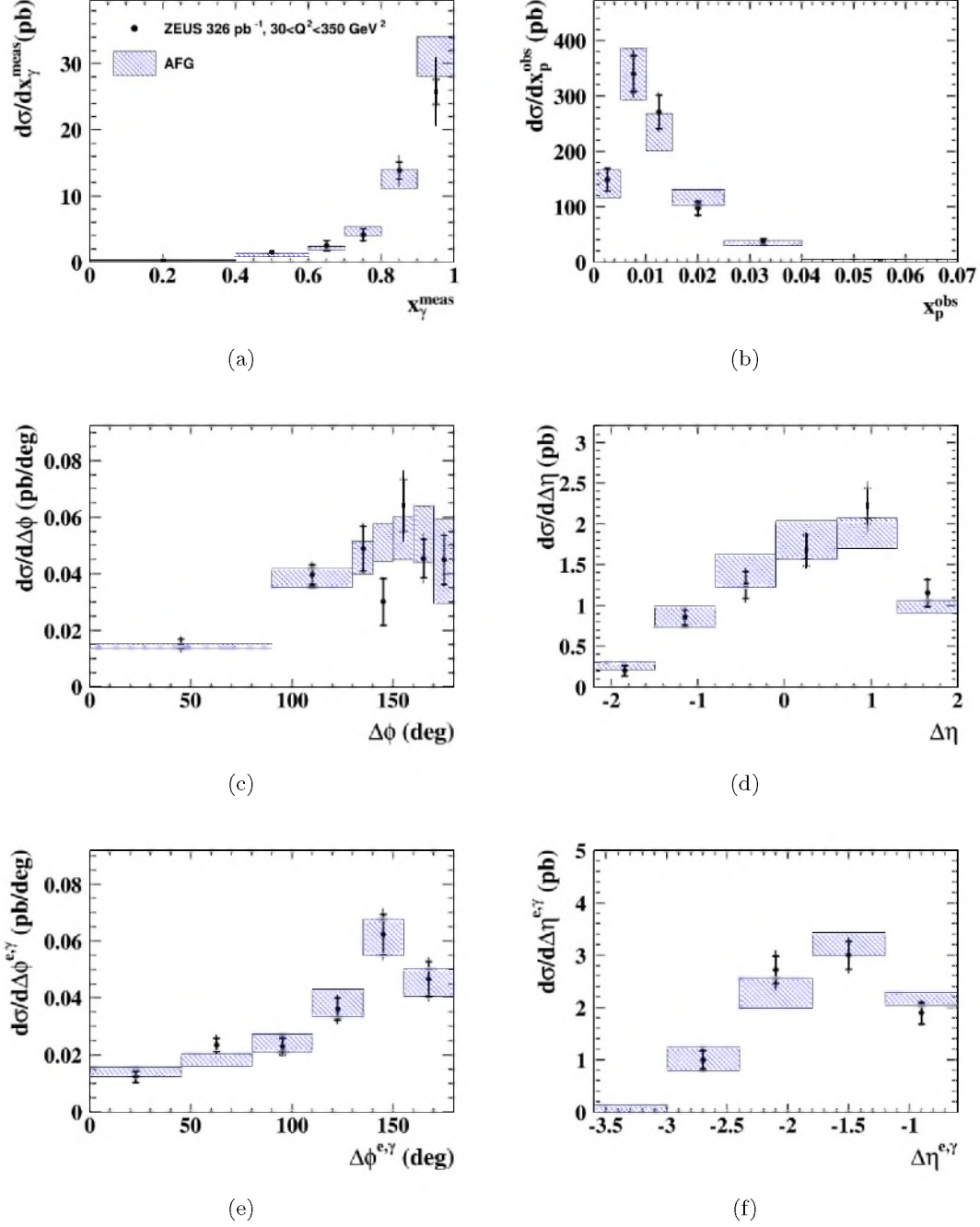


Figure 9. Differential cross sections for selected variables in the region $30 \leq Q^2 < 350 \text{ GeV}^2$ as in figures 4, 5. Theoretical predictions from Aurenche et al. (AFG) are shown, with associated uncertainties indicated by the bands.

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